

TOPIC \div Electronic partition functions

Pages (9-10)

AT STP, i.e., $T = 298\text{ K}$

Dr. P. K. Hazra

$$\therefore k_B T \approx 9.43 \times 10^{-4} \text{ a.u.}$$

$$\approx 9.43 \times 10^{-4} \text{ hartree}$$

\Rightarrow Electronic levels appear
discretized for thermal energy
&
exchange of energy at
thermal levels

Thus; electronic partition functions, q_{elec} can be represented as \sum_i

$$q_{elec} = \sum_i w_i e^{-\beta \epsilon_i}$$

$\forall w_i =$ degeneracy of energy level " ϵ_i "

Since all thermodynamic properties constitute equations with q_{elec} both in numerator and denominator and because of negative bound states of atomic & molecular systems, ϵ_0 can be set to zero. Thus

$$q_{elec} = w_{e1} + w_{e2} e^{-\beta \Delta \epsilon_{12}}$$

Usually, Not more than one or two excited states contribute to " q_{elec} ".

TOPIC 2

Nuclear Partition Function

For transition among nuclear energy levels, required temperature, $T \approx 10^{10}$ K, Thus at STP, majorly ground states involved.

$$\Rightarrow q_{nuc} = \sum W_{nj} e^{-\beta \epsilon_{nj}}$$

Can be written as

$$q_{nuclear} = W_{n1} + \dots$$

Higher levels can be ignored.

As Born-Oppenheimer approximation separates electronic & nuclear coordinates, so does their respective partition functions

Hence total partition function of monoatomic gases can be

written as

$$Z_{monoatomic} = \frac{(q_{trans} q_{elec} q_{nuc})^N}{N!}$$

where

$$q_{trans} = \left(\frac{2\pi m k_B T}{h^2} \right)^{3/2} V = \frac{V}{\Lambda^3}$$

$$q_{elec} = W_{e1} + W_{e2} e^{-\beta \Delta \epsilon_{12}}$$

$$q_{nuc} = W_{n1} + \dots$$