Principles of Monte Carlo Simulation

What is Monte Carlo Simulation?

- MC is a stochastic simulation technique based on random numbers and probability.
- Widely used across various disciplines in Economics, Physics, Chemistry, Biophysics and Material Science.
- Typically used to solve complex many-body problems with different types of interactions.
- Classical Monte Carlo: Uses Boltzmann distribution for random sampling to calculate the thermodynamic properties
- Quantum Monte Carlo: Random walks are used to calculate energies and wavefunctions using Schrodinger's equation.

Why Monte Carlo Simulation?

- It is used for problems which are difficult to solve mathematically.
- Too many particles in the system.
- Complex interactions among particles and/or in presence of an external field.

Properties of MC Simulation

- Large number of inputs.
- The inputs should be random and uniformly distributed.

An Example of MC Simulation

Value of π may be determined by MC simulation.

Inscribe a circle within a square.

Some trial shots are generated uniformly over the square.

The total number of trial shots inside the shaded area of the circle is N_{hit} and the total number of shots are N_{shot}

$$\pi \approx \frac{4N_{hit}}{N_{shot}}$$



Calculation of Definite Integrals by Conventional Methods



n uniformly separated points

Calculation of Definite Integrals by Monte Carlo Method



Classical Monte Carlo Simulation

$$\langle A \rangle = \frac{1}{Z} \iint dp \, dq \, A(p,q) \exp\{-\beta H(p,q)\}$$

where Z is defined as $Z = \frac{1}{h^{3N}} \iint dp \, dq \exp\{-\beta H(p,q)\}$, H(p,q) is the Hamiltonian that describes the system and $\beta = \frac{1}{k_BT}$.

An exact calculation of the partition function Z is not possible since it involves an integration over 6N dimensional phase space.

Calculation of the complex many-body interaction in the Hamiltonian is also difficult.

Monte Carlo Sampling

- It is difficult to compute the *6N* dimensional integral for computing the ensemble average.
- Therefore average is calculated over a finite number of states, a subset of all possible states, with the correct Boltzmann weight. If a subset is chosen at random with a probability distribution p_x then

$$A_N = \frac{\sum_{\chi=1}^N A_{\chi} p_{\chi}^{-1} \exp(-\beta E_{\chi})}{\sum_{\chi=1}^N p_{\chi}^{-1} \exp(-\beta E_{\chi})} \qquad N \to \infty, \ A_N \approx < A >$$

Usually in MC simulation N is very large i.e., $10^6 - 10^{10}$

Simple Sampling

- To specify p_x such that the chosen N states will give correct < A >
- A choice of $p_X = p$ gives

$$A_{N} = \frac{\sum_{x=1}^{N} A_{x} p^{-1} \exp(-\beta E_{x})}{\sum_{x=1}^{N} p^{-1} \exp(-\beta E_{x})} = \frac{\sum_{x=1}^{N} A_{x} \exp(-\beta E_{x})}{\sum_{x=1}^{N} \exp(-\beta E_{x})}$$

This average is valid for non-thermal and non-interacting systems of molecules.

For such systems, the temperature, T, is high

$$\langle A \rangle \approx A_N = \frac{1}{N} \sum_{x=1}^N A_x$$

Importance Sampling

• For discrete energy states, the ensemble average of A is given by

$$\langle A \rangle = \frac{1}{Z} \sum_{x=1}^{N} A_x \exp(-\beta E_x)$$
 where $Z = \sum_{x=1}^{N} \exp(-\beta E_x)$

- At high temperature, the Boltzmann factor is unity which reduces to simple average.
- At low temperature, the system is in the ground state or one of the lowest energy states.
- A small fraction of the phase space only contributes to the average and hence this part of the phase space should be scanned.
- Random sampling of the phase space is therefore computationally inefficient.

Metropolis Monte Carlo

- An atom is chosen randomly and is subjected to an uniform random displacement along each of the spatial coordinates, *x*, *y* and *z*.
- Random number generator is used to generate the new position coordinate. This is a FORTRAN program that gives uniformly distributed random numbers between [0,1].
- The size of the move is controlled by δr_{max} .
- The transition between two states depend on the relative probabilities of the initial state *m* and the final state *n* which is given by

$$P_m = \frac{\exp(-\beta E_m)}{Z}$$
 and $P_n = \frac{\exp(-\beta E_n)}{Z}$

- The probabilities P_m and P_n cannot be determined exactly as it is difficult to calculate the canonical partition function.
- The ratio of the probabilities of the initial and final states is $\frac{P_m}{P_n} = \exp(-\beta E_{mn}) \quad \text{where} \quad E_{mn} = E_n - E_m$ • Transition probability = $\frac{P_m}{P_n} = \exp(-\beta E_{mn}) \quad if \ E_{mn} > 0$

$$\frac{P_m}{P_n} = 1 \quad if \ E_{mn} \le 0$$

Metropolis MC Method

Call a random number, r = [0,1] uniformly distributed between 0 and 1.

If $r \leq \frac{P_m}{P_n}$, there is a change in state from *m* to *n*, otherwise the initial state is retained.







Application of Metropolis MC

Different conditions for the acceptance/rejection of moves

If $E_n < E_m$

The MC move is accepted and E_n becomes the energy of new conformation.

If Transition probability > *r*

Then the MC move is accepted and E_n becomes the energy of the new conformation. If $E_n > E_m$

The MC move is accepted conditionally

Transition probability is calculated

and compared with a random number *r* between 0 and 1.

If Transition probability < r

Then the MC move is rejected and the energy of the initial state E_m is retained.

Choice of Potential

- Determination of the potential energy is an important step in MC and MD simulations.
- The change in potential energy is given by $\delta E_{mn} = \sum_{j=1}^{N} E_{ij}^n - \sum_{j=1}^{N} E_{ij}^m$

Sum over atoms excludes the *i*-th atom.

• The interaction of the atom i with its neighbours up to a certain cutoff distance R_c is considered.

Minimum Image Convention

- Potential energy is calculated using the pairwise additivity principle.
- For a short-ranged potential energy function, the interaction between an atom with other atoms and its closest periodic images that are within a cut-off distance R_C are only considered.
- Natural consequence of the periodic boundary condition.



Potential Energy Plot



Density Plot

